

## A. A Formal and Exact Approach

- Turn TISE  $\hat{H}\psi = E\psi$  into a huge Matrix problem

What for?

- Finding eigenvalues / eigenvectors of a huge matrix is a standard problem of computational physics [numerical package]
- Starting point for developing / understanding approximations

$\hat{H}\psi = E\psi$  (A1) Often, know the equation but can't  
solved it analytically (解析解)

- But, still possible to solve TISE exactly (精确解) or almost exactly by numerical approaches (computationally)
- Inspect QM problem and domain (region of space)
- Choose a (convenient) set of "basis functions"  
 $\{\phi_1, \phi_2, \phi_3, \dots, \phi_i, \phi_j, \dots\}$  (often infinitely many)  
 such that any other function (especially eigenstates  $\psi$  in Eq.(1))  
 can be expanded in terms of  $\underbrace{\{\phi_i\}}$   
 Symbol means the whole set



- Functions  $\{\phi_i\}$  can always be chosen to be normalized
- For generality, NOT assuming  $\{\phi_i\}$  to be orthogonal  
 [Why not? Can't do it for some TISE problems, e.g. molecules]

See Eq. (3), there are  $\int \phi_j^* \phi_i d\tau = \langle \phi_j | \phi_i \rangle$  (Dirac notation)



- take  $\phi_j$  &  $\phi_i$ , form an integral
- it is a number!

- This number can be labelled by  $j$  and  $i$

Call  $\int \phi_j^* \phi_i d\tau \equiv S_{ji}$  (A4)

$S_{ii} = \int \phi_i^* \phi_i d\tau = 1$  (normalized), but  $S_{ji} \neq 0$  in general for  $i \neq j$

There are also  $\int \phi_j^* \hat{H} \phi_i d\tau = \langle \phi_j | \hat{H} | \phi_i \rangle$

a quantity of unit "energy" labelled by  $j$  and  $i$

Call  $\int \phi_j^* \hat{H} \phi_i d\tau \equiv H_{ji}$  (know  $\hat{H}$  and  $\{\phi_i\}$ , then let the computer do the integrals  $H_{ji}$ )

∴ Eq. (3) becomes

$\sum_i H_{ji} a_i = E \sum_j S_{ji} a_i$  OR  $\sum_i (H_{ji} - E S_{ji}) a_i = 0$  (A6)

equivalent to TISE (Eq. (A1)) (等價)

- Eq. (A1) is an energy eigenvalue problem
- Schrödinger form: TISE is a differential equation
- Equivalent form in Eq. (A6): turned TISE into a Matrix problem

Refresher: Matrix  $\vec{M}$  multiplies into a column vector

$$\underbrace{\begin{pmatrix} M_{11} & M_{12} \\ M_{21} & M_{22} \end{pmatrix}}_{\substack{\textcircled{2} \times \textcircled{2} \\ \uparrow \text{same} \\ \textcircled{2} \times \textcircled{1}}} \underbrace{\begin{pmatrix} a_1 \\ a_2 \end{pmatrix}}_{\textcircled{2} \times \textcircled{1}} = \underbrace{\begin{pmatrix} M_{11}a_1 + M_{12}a_2 \\ M_{21}a_1 + M_{22}a_2 \end{pmatrix}}_{\textcircled{2} \times \textcircled{1}}$$

Inspect:  $M_{11}a_1 + M_{12}a_2 = \sum_{i=1,2} M_{1i}a_i$  — (i) (first row)

$$M_{21}a_1 + M_{22}a_2 = \sum_{i=1,2} M_{2i}a_i$$
 — (ii) (second row)

So,  $j=1$ ,  $\sum_{i=1,2} M_{ji}a_i$  is (i)  $\searrow$   
 $j=2$ ,  $\sum_{i=1,2} M_{ji}a_i$  is (ii)  $\swarrow$   $\therefore \sum_{i=1,2} M_{ji}a_i$  (for various  $j$ )  
 is a matrix multiplying into a column vector

Back to Eq. (A6):  $\sum_i (H_{ji} - ES_{ji}) a_i = 0$

It is... 
$$\begin{pmatrix} H_{11} - ES_{11} & H_{12} - ES_{12} & H_{13} - ES_{13} & \dots & \dots \\ H_{21} - ES_{21} & H_{22} - ES_{22} & H_{23} - ES_{23} & \dots & \dots \\ H_{31} - ES_{31} & H_{32} - ES_{32} & H_{33} - ES_{33} & \dots & \dots \\ \vdots & \vdots & \vdots & \ddots & \vdots \\ \vdots & \vdots & \vdots & \ddots & \vdots \end{pmatrix} \begin{pmatrix} a_1 \\ a_2 \\ a_3 \\ \vdots \\ \vdots \end{pmatrix} = 0$$

Key Result  $\rightarrow$  (A7)  
(same as (A6))

- General so far, leave  $S_{ji}$  general, equivalent to TISE
- Typically,  $\begin{pmatrix} \dots \\ \dots \\ \dots \end{pmatrix}$  is  $\infty \times \infty$  Matrix ( $\because$  Infinitely many  $\phi_i$ 's)
- TISE unknowns are  $E$  and  $\psi$  (many pairs)
- Eq. (7):  $E$  and  $\begin{pmatrix} a_1 \\ a_2 \\ \vdots \end{pmatrix}$  are unknowns (many pairs)

▪  $\begin{pmatrix} 0 \\ 0 \\ 0 \\ \vdots \end{pmatrix}$  is obviously a "solution", but  $\psi = \sum_i a_i \phi_i = 0$  (particle disappears!)  
 trivial solution

▪ For non-trivial solutions, requires determinant

$$\begin{vmatrix} H_{11} - ES_{11} & H_{12} - ES_{12} & H_{13} - ES_{13} & \dots \\ H_{21} - ES_{21} & H_{22} - ES_{22} & H_{23} - ES_{23} & \dots \\ H_{31} - ES_{31} & H_{32} - ES_{32} & H_{33} - ES_{33} & \dots \\ \vdots & \vdots & \vdots & \ddots \end{vmatrix} = 0 \quad (\text{A8})$$

- Plug in a value of  $E$ , LHS gives a value, but may not be zero  
 $\Rightarrow$  Eq. (A8) is not satisfied  $\Rightarrow$  wrong  $E$  (then try again...)
- ∴ Eq. (A8) is an equation for allowed values of  $E$   
 $|\infty \times \infty| = 0 \Rightarrow$  infinite many allowed  $E$   
 not surprising (c.f. 1D well, harmonic oscillator, ...)



- Practically, truncate into  $N \times N$  size (approximation), then  $|N \times N| = 0 \Rightarrow N$  values of allowed energy

### Take-home Message

- Choose a set of  $\{\phi_i\}$ , TISE becomes Eq. (A7) in Matrix form
- Completely general up to here, thus Eq. (A7) is exact
- Cleverer choice(s) of  $\{\phi_i\}$  can shorten computing time
- Take Eq. (A7) in, we shall use it many times

### Further questions

- Is  $H_{ji}$  related to  $H_{ij}$ ? If yes, how? [ $\hat{H}$  is a Hermitian operator]
- How about  $S_{ji}$  and  $S_{ij}$ ?

## Refresher: Hermitian Operators

$$\int f^* \hat{A} g \, d\tau = \int g \hat{A}^* f^* \, d\tau = \left( \int g^* \hat{A} f \, d\tau \right)^* \quad \text{for Hermitian } \hat{A}$$

↑ any  $f$  and  $g$

- Hamiltonian  $\hat{H}$  is Hermitian (allowed energies are real)

$$\int \phi_j^* \hat{H} \phi_i \, d\tau = \left( \int \phi_i^* \hat{H} \phi_j \, d\tau \right)^*$$

$$H_{ji} = H_{ij}^* \quad (\hat{A}^\dagger)$$

Matrices obeying this relation are called Hermitian Matrices

( $ji$ )-element = complex conjugate of ( $ij$ )-element

$\therefore$  Only need to calculate half of matrix elements

and  $H_{ii} = H_{ii}^* \Rightarrow$   $H_{ii}$  are real (diagonal elements are real)

Similarly,  $S_{ji} = S_{ij}^*$  ( $\because$  " $\hat{1}$ " is Hermitian)

### Simplified Form of Eq. (A7)

If  $\{\phi_i\}$  give  $\int \phi_j^* \phi_i d\tau = 0$  or  $\approx 0$ ,  
orthogonal  
for  $j \neq i$   
(approximately orthogonal)

and  $\int \phi_i^* \phi_i d\tau = 1$ , then Eq. (6) becomes

$$\sum_j H_{ji} a_i = E \sum_j \underbrace{S_{ji}}_{\delta_{ji}} a_i = E a_j \quad (A10)$$

$$\begin{pmatrix} H_{11} & H_{12} & H_{13} & \dots \\ H_{21} & H_{22} & H_{23} & \dots \\ H_{31} & H_{32} & H_{33} & \dots \\ \vdots & \vdots & \vdots & \ddots \\ \vdots & \vdots & \vdots & \ddots \end{pmatrix} \begin{pmatrix} a_1 \\ a_2 \\ a_3 \\ \vdots \\ \vdots \end{pmatrix} = E \begin{pmatrix} a_1 \\ a_2 \\ a_3 \\ \vdots \\ \vdots \end{pmatrix} \quad (A11)^+$$

which is clearly an eigenvalue problem of the matrix  $H_{ji}$

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<sup>+</sup> Eq. (A11) can be obtained directly from Eq. (A7) by setting  $S_{ii} = 1$  &  $S_{ji} = 0$  ( $j \neq i$ )

## Further Remarks

- Eq. (A7) or Eq. (A11) (if valid) is the starting point in understanding electronic states in molecules (molecular orbitals) and in solids (energy bands)
- Starting point for computational approaches
- Truncation: Want lowest 10 eigenvalues  $E$ , truncate to  $500 \times 500$  (an approximation) may not affect the lowest eigenvalues
- If Eq. (A11) is valid, allowed energies  $E$  are eigenvalues of

$$\overleftrightarrow{H} = \begin{pmatrix} H_{11} & H_{12} & H_{13} & \dots \\ H_{21} & H_{22} & H_{23} & \dots \\ H_{31} & H_{32} & H_{33} & \dots \\ \vdots & \vdots & \vdots & \ddots \end{pmatrix} \quad (A12)$$

- QM TISE problems are related to Matrix Mathematics (Heisenberg/Born)

• Need a basis  $\{\phi_i\}$  to construct the matrix  $\hat{H}$  (Recall)

• Keep in mind that

Good

$$\begin{array}{c}
 \langle \phi_1 | \\
 \langle \phi_2 | \\
 \langle \phi_3 | \\
 \vdots \\
 \vdots
 \end{array}
 \begin{array}{c}
 \underline{|\phi_1\rangle} \quad \underline{|\phi_2\rangle} \quad \underline{|\phi_3\rangle} \quad \dots \\
 \left( \begin{array}{cccc}
 H_{11} & H_{12} & H_{13} & \dots \\
 H_{21} & H_{22} & H_{23} & \dots \\
 H_{31} & H_{32} & H_{33} & \dots \\
 \vdots & \vdots & \vdots & \vdots \\
 \vdots & \vdots & \vdots & \vdots
 \end{array} \right)
 \end{array}$$

there is a basis set

and

$$\begin{array}{c}
 a_1 \\
 a_2 \\
 a_3 \\
 \vdots \\
 \vdots \\
 \vdots
 \end{array}
 \begin{array}{c}
 |\phi_1\rangle \\
 |\phi_2\rangle \\
 |\phi_3\rangle \\
 \vdots \\
 \vdots \\
 \vdots
 \end{array}
 \quad (A13)$$

$H_{ij}$  depends on  $\hat{H}$  AND  $\{\phi_i\}$

meaning

$$\psi = a_1 \phi_1 + a_2 \phi_2 + a_3 \phi_3 + \dots$$

many choices [some more convenient or obvious]

"Cleverest" Choice

$$\hat{H} \psi_i = E_i \psi_i$$

↑ energy eigenstates

Choose the basis set to be  $\{\psi_i\}$ , then  $\int \psi_j^* \hat{H} \psi_i d\tau = E_i \int \psi_j^* \psi_i d\tau$

$E_i \psi_i$   
orthonormal

Meaning:  $H_{ii} = E_i$ ;  $H_{ij} = 0$  ( $i \neq j$ )

$= E_i \delta_{ij}$

$$\therefore \vec{H} = \begin{pmatrix} E_1 & & & \\ & E_2 & & \\ & & E_3 & \\ & & & \ddots \end{pmatrix}$$

is diagonal & elements are eigenvalues

- If Eq. (A11) is valid, solving Eq. (A11) is equivalent to "diagonalizing the matrix  $\vec{H}$ " in (A12), i.e. changing (looking for) the basis from  $\{\phi_i\}$  ( $\vec{H}$  not diagonalize) to  $\{\psi_i\}$  ( $\vec{H}$  is diagonalized)

- Finally, if  $\infty \times \infty$  or  $N \times N$  is scary, don't worry.
- Get yourself familiar with a simple case
  - focus on two functions  $\phi_i$  and  $\phi_j$  (forget the others)

So only  $H_{ii}$ ,  $H_{ij}$ ,  $H_{ji}$ ,  $H_{jj}$

$$\begin{pmatrix} H_{ii} - ES_{ii} & H_{ij} - ES_{ij} \\ H_{ji} - ES_{ji} & H_{jj} - ES_{jj} \end{pmatrix} \quad \text{in Eq. (A7)}$$

and

$$\begin{pmatrix} H_{ii} - E & H_{ij} \\ H_{ji} & H_{jj} - E \end{pmatrix} \quad \text{in Eq. (A11)}$$

↳ or eigenvalue problem of  $\begin{pmatrix} H_{ii} & H_{ij} \\ H_{ji} & H_{jj} \end{pmatrix}$

∴ Knowing how to treat  $2 \times 2$  matrices is useful in many QM problems  
 [Any everyone can handle  $2 \times 2$  matrices!]